

PicoTesla absolute field readings with a hybrid $^3\text{He}/^{87}\text{Rb}$ magnetometer

Christopher Abel^{1,a}, Georg Bison^{2,b}, W. Clark Griffith¹, Werner Heil³, Klaus Kirch^{2,4}, Hans-Christian Koch², Bernhard Lauss², Alexander Mtchedlishvili², Martin Pototschnig⁴, Philipp Schmidt-Wellenburg², Allard Schnabel⁵, Duarte Pais², and Jens Voigt⁵

¹ Department of Physics and Astronomy, University of Sussex, Brighton, UK

² Paul Scherrer Institute, 5232 Villigen, Switzerland

³ Institut für Physik, Johannes-Gutenberg-Universität, 55128 Mainz, Germany

⁴ Institute for Particle Physics and Astrophysics, ETH Zurich, 8093 Zurich, Switzerland

⁵ Physikalisch-Technische Bundesanstalt, Berlin, Germany

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Abstract. We demonstrate the use of a hybrid $^3\text{He}/^{87}\text{Rb}$ magnetometer to measure absolute magnetic fields in the pT range. The measurements were undertaken by probing time-dependent ^3He magnetisation using ^{87}Rb zero-field magnetometers. Measurements were taken to demonstrate the use of the magnetometer in cancelling residual fields within a magnetic shield. It was shown that the absolute field could be reduced to the 10 pT level by using field readings from the magnetometer. Furthermore, the hybrid magnetometer was shown to be applicable for the reduction of gradient fields by optimising the effective ^3He T_2 time. This procedure represents a convenient and consistent way to provide a near zero magnetic field environment which can be potentially used as a base for generating desired magnetic field configurations for use in precision measurements.

1 Introduction

Many precision measurements require a thorough knowledge of surrounding magnetic fields, for example fundamental physics searches [1–3], electric dipole moments [4–7], exotic spin coupling forces [8,9], and axions [10–13], as well as general calibration of magnetic sensors [14]. It is generally difficult to produce an absolutely correct magnetic field with precise knowledge of both the direction and magnitude, especially due to the inevitable presence of small residual fields after typical degaussing procedures of magnetic shielding. However, if one starts with an environment where the residual magnetic field is as close to zero as possible, any field can then in principle be generated on top of this using one or more calibrated coils.

We present here a hybrid $^{87}\text{Rb}/^3\text{He}$ magnetometer that can be used to provide absolute magnetic field readings based on the free spin precession (FSP) of ^3He , monitored by optical ^{87}Rb magnetometers that detect the small magnetic field variations generated by the precessing ^3He polarisation. This magnetometer can be used to zero a residual field environment, upon which a desired magnetic field can be applied using a separate coil system.

This work follows on from a previously developed combined $^3\text{He}/^{133}\text{Cs}$ magnetometer described in [15], capable of measuring magnetic fields around 1 μT with sensitivities of ~ 50 fT per 180 s measurement cycle. In this work the ^{133}Cs readout sensors based on rf-driven magnetic resonance are replaced with ^{87}Rb sensors capable of operating at zero magnetic field. A similar concept was demonstrated by Cohen-Tannoudji et al. in 1969 [16], where a low-field $^3\text{He}/^{87}\text{Rb}$ magnetometer was used to measure an applied 200 pT holding field. Here we apply no explicit holding field but instead allow the ^3He to precess around the residual field inside a degaussed magnetic shield. After applying a field minimisation process with a set of correction coils, we have measured absolute fields below 5 pT. We have also demonstrated minimisation of field gradients at the ^3He cell by maximising the effective T_2 relaxation time.

2 Measurement principle

Measuring low magnetic fields with a ^3He FSP magnetometer follows two steps; polarisation of ^3He in an enclosed volume, followed by detection of the free precession signal.

Metastability exchange optical pumping (MEOP) was used to polarise ^3He which has been shown to achieve

^a e-mail: c.abel@sussex.ac.uk

^b e-mail: georg.bison@psi.ch

polarisation levels up to $p \sim 80\%$ [17]. At sufficiently low fields, when the Larmor frequency of the ^3He is much less than the polarisation pump rate, the helium can be polarised without the need for a holding field, such as that used in [16]. During MEOP, the ^3He polarisation is held in the direction of the pump laser propagation \mathbf{k}_{He} . After polarisation, when the laser radiation is switched off, the ^3He polarisation precesses around the volume-averaged cell field vector. The frequency of the ^3He free-precession is free from light shifts commonly associated with optically pumped magnetometers [18]. The precessing magnetisation of the ^3He is then probed using zero-field ^{87}Rb magnetometers (RbZFM) placed close to the ^3He cell. Using the Larmor frequency of the ^3He , the volume-averaged field magnitude at the centre of the cell can be calculated along with values for the relaxation time of the polarisation (T_2). The high sensitivity of the ^{87}Rb zero field magnetometer to field changes of well below 1 pT (AC below 100 Hz) is used to precisely detect perturbations caused by the precessing ^3He magnetisation at the ^{87}Rb magnetometer's position.

3 Experimental apparatus

The hybrid magnetometer is schematically shown in Figure 1, and a photograph in Figure 2. Two commercially available ^{87}Rb zero-field magnetometers¹ were used for the measurements to observe the spin precession of the ^3He inside the glass cell. Each RbZFM unit contains a cubic rubidium cell with side 6 mm, through which a 795 nm laser tuned to the D1 line of ^{87}Rb is passed for optical pumping and detection [19]. A 923 Hz oscillating magnetic field is applied to the atoms in the cell. Two separate orthogonal magnetic fields (with a 90° phase shift) are applied across the cell, which, when demodulated, enables two RbZFM channels to be measured per sensor-head. The magnetometer has a dynamic range of ± 5 nT. Use of the sensor in two-axis mode reduces the sensitivity by around 30%. The internal correction coils of the ^{87}Rb magnetometers were not used, as the local DC fields affect the residual field value at the ^3He cell.

In this experiment a 70 mm diameter 1 mbar ^3He cell was used. If the ^3He was 100% polarised, then a ^{87}Rb cell located 39 mm from the centre of the ^3He cell, orientated 45° from the initial polarisation direction, would measure a field of 80 pT along the Z axis. The ^3He was polarised by optical pumping of the gas using circularly polarised 1083 nm light from a commercially available 5 W ytterbium doped fibre laser². The RbZFMs were orientated to enable at least one reading in each of three axes. The hybrid magnetometer was placed in a 5-layer μ -metal shield, allowing external magnetic noise to be suppressed by a factor $\sim 10^5$. The shield housed six coils enabling independent control for all three field components, and first-order gradients. The ^3He was polarised using the MEOP process, during which a discharge is struck inside the cell. After polarisation the discharge is switched off and the

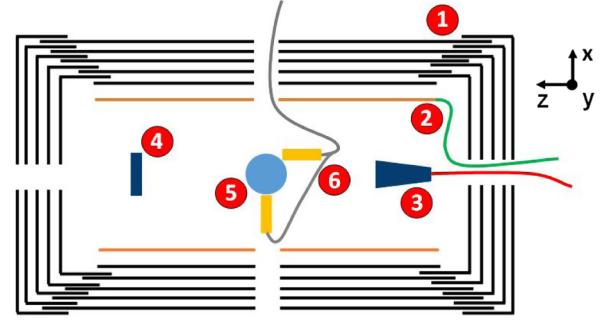


Fig. 1. Schematic diagram of the hybrid magnetometer, along with the axis definitions. (1) magnetic shielding, (2) internal coils, (3) helium laser expansion optics, (4) 1083 nm back-reflecting mirror, (5) ^3He cell, (6) ^{87}Rb magnetometers. The outermost shield layer has a length 200 cm and a diameter 94 cm. The innermost layer has length 194 cm and diameter 50 cm.

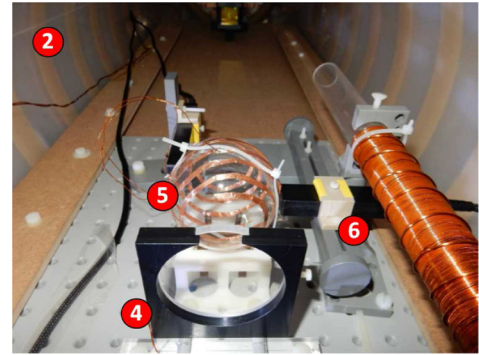


Fig. 2. Photograph of the ^{87}Rb and ^3He magnetometers being installed into the shield with elements labelled according to Figure 1. The helium laser expansion optics are just visible at the top of the image. The Tesla coil on the right was used to ignite a discharge in the ^3He cell to enable MEOP.

laser ramped down before turning off. After the ^3He has been polarised the ^{87}Rb magnetometers probe the precessing ^3He dipole.

The entire experiment was located in a closed area with a controlled temperature stability of $\pm 1^\circ\text{C}$. A data acquisition system recorded the outputs from the RbZFM electronics module with a 200 Hz sampling rate. The data collection device was interfaced with a constant current supply, simultaneously enabling control of the current in three correction coils.

4 Field minimisation

4.1 Method

The spin precession frequency f_{He} can be converted to an absolute field at the ^3He cell position by

$$|B_{\text{He}}| = f_{\text{He}} \frac{2\pi}{\gamma_{\text{He}}}, \quad (1)$$

¹ QuSpin QZFM [19].

² Keopsys CYFL-GIGA.

where $\gamma_{\text{He}}/2\pi = 32.43410084(81) \text{ Hz}/\mu\text{T}$ [20] is the gyro-magnetic ratio of ^3He . The field measured by the ^3He cell inside the μ -metal shield is the absolute value of the residual field B_R , which can be expressed by the orthogonal axis components according to

$$|B_R| = \sqrt{B_{Rx}^2 + B_{Ry}^2 + B_{Rz}^2}. \quad (2)$$

The shield was equipped with coils which produce additional fields along the different axes. Assuming perfectly orientated coils and neglecting inhomogeneities of the produced fields, the residual field together with the coil contribution can be expressed as

$$|B_{R+C}|(I_x, I_y, I_z) = \sqrt{(\mathcal{B}_x)^2 + (\mathcal{B}_y)^2 + (\mathcal{B}_z)^2}, \quad (3)$$

where $\mathcal{B}_i = B_{Ri} - A_i I_i$, with A and I the coil coefficients and the currents applied to the coils in the orthogonal directions, respectively, and $i = x, y, z$. Each summand under the square root can be set to zero independently by adjusting the corresponding current; this can easily be done by minimising the function $|B_{R+C}|(I_x, I_y, I_z)$ sequentially. After adjusting I_x , I_y and I_z , $|B_{R+C}|$ should be zero if the coils are perfectly aligned and the ^3He cell is small in diameter so that gradient effects can be neglected. As long as equation (3) is valid, and if the coil coefficients A_x , A_y and A_z are measured in advance, the components of the residual field can be calculated by

$$B_{Ri} = A_i I_{i,\text{optimum}}. \quad (4)$$

For the more general case including some misalignment or inhomogeneity in the fields generated by the coils, equation (3) can be written as

$$|B_{R+C}| = \sqrt{\sum_{i=1}^3 \left(B_{Ri} - \sum_{j=1}^3 A_{ij} I_j \right)^2}, \quad (5)$$

where A_{ij} is now a matrix with A_x now denoted as A_{xx} , A_y as A_{yy} , and A_z as A_{zz} . As long as the field vectors generated by different coils are not linearly dependent, the matrix A_{ij} should be invertible, and there will be exactly one set of currents that will compensate a particular residual field. Equation (5) can be minimised by sequentially adjusting each coil current to find the minimum field point, which if the vectors that make up A_{ij} are not mutually orthogonal will not reach zero field after a single iteration of the three currents. Further iterations should allow further reduction in $|B_{R+C}|$, and in principle should eventually converge to zero field, although in practice this could be very time consuming. Alternatively, if the full A_{ij} matrix can be determined in advance then this might allow a more clever and quicker optimisation method to be found.

The accuracy of the measurements strongly depends on the ability to determine the frequency of the spin precession from the RbZFM signals. In the experiment, a mean frequency reading was taken from all the RbZFM channels with an appropriate signal-to-noise ratio. With

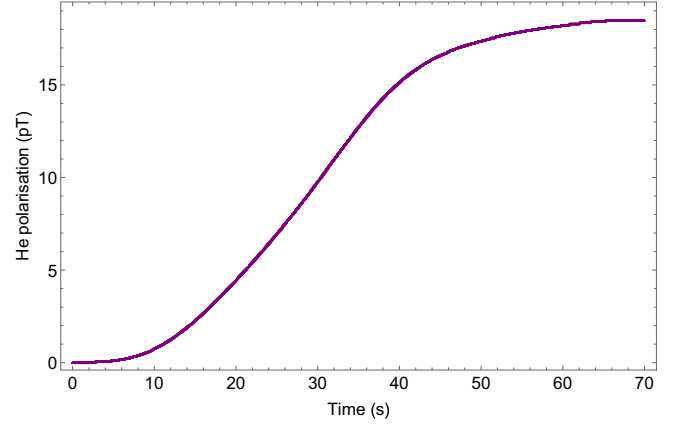


Fig. 3. A plot displaying the field in the Z direction as measured by a ^{87}Rb magnetometer during the ^3He polarisation procedure. The location of the ^{87}Rb magnetometer corresponds to the geometry described in Section 3, from which the ^3He polarisation can be approximated at 24%.

two RbZFM devices, each containing two channels perpendicular to each other, one field direction is measured twice. When precessing, the ^3He magnetisation precesses around B_R . If the initial polarisation (\mathbf{k}_{He}) happens to be in the same direction as B_R , no precession will occur. The RbZFM devices are least sensitive along the direction of B_R , resulting in the hybrid magnetometer being arranged to measure three orthogonal axes. If the field components are optimised sequentially the Z component should be optimised first as the signal amplitude approaches zero when $\mathbf{B}_R \rightarrow (0, 0, B_R)$. In future measurements it would be advantageous to place the two RbZFM sensor-heads at $+X$ and $-X$, orientated in such a way that the orthogonal channels can also measure Y and Z. In this orientation the two opposing X channels can produce a gradiometer signal, improving the signal-to-noise ratio of the ^3He precession by the reduction of common-mode noise [21].

The principal limitation to the measurements is the long precession period of ^3He at near-zero fields. Although the measurements are fundamentally limited by the finite T_2 time of the ^3He polarisation, if we limit our measurement time for one reading to 1 h, the absolute value of the residual field cannot be lower than 10 pT according to equation (1). In the time it takes to record a single precession at these field values, other systematic effects can dominate the reading, such as the time stability of the RbZFM readings and the field inside the shield. Also, as the Earth's frame rotates around the freely-precessing ^3He at a rate of an inverse sidereal day, the ^{87}Rb magnetometers will measure a shift in frequency up to $|\Delta f_{\text{He}}| \approx 12 \mu\text{Hz}$ resulting in $|\Delta B_R| \approx 357 \text{ fT}$.

4.2 Results and discussion

All data were taken during May 2018 at the Paul Scherrer Institute (PSI) in Villigen, Switzerland.

Figure 3 shows the magnetisation build up in the ^3He cell as a function of time in the magnetic shielding, as

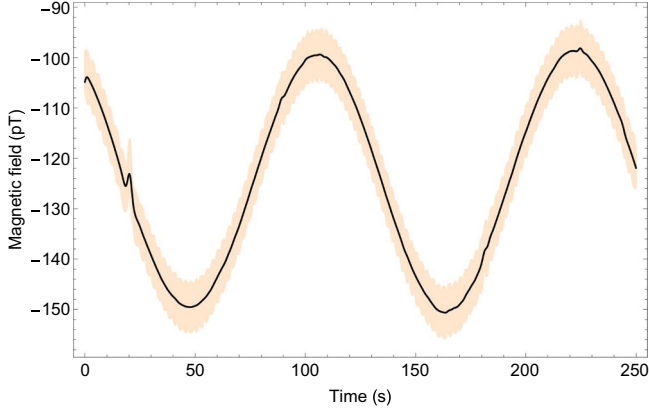


Fig. 4. Raw ^{87}Rb data displaying the oscillating ^3He signal and 1 Hz low-pass filtered data (black). The field shifts at 25 and 230 s were caused by small mechanical disturbances of the magnetic shielding. The precession frequency corresponds to a field reading of 264.9 pT.

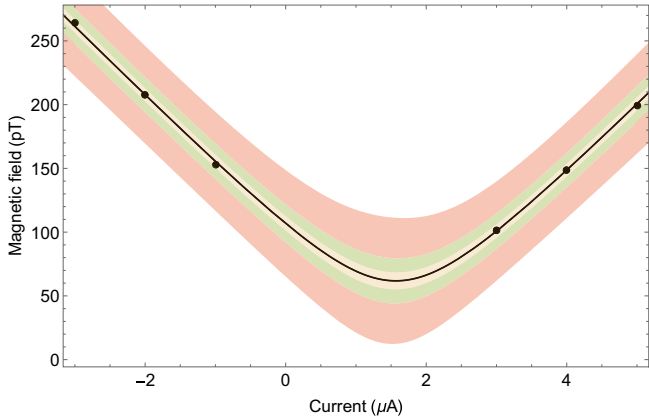


Fig. 5. Scan along X fitted with equation (3). The bands correspond to the 1, 2 and 3 σ levels. The scan was performed after the Y component of the field had already been minimised.

measured by the ^{87}Rb magnetometers. We first measured the magnitude of the post-degauss residual field inside the shield, a plot of which can be seen in Figure 4. By using the gyromagnetic ratio of ^3He the fitted frequency corresponds to a residual magnetic field of 264.9 pT. The raw data were filtered via a 1 Hz low-pass filter, then a sine fit applied to the filtered data.

Data were taken to determine the effectiveness of the minimisation technique in reducing residual fields inside the shield, after the field had been minimised along all three field components. After the three corresponding field currents were applied to the correction coils, the true field magnitude at the predicted minima was recorded, thereby offering a comparison to the predicted minimum field from the fitting routine, and the field magnitude proper.

Figures 5 and 6 show the results of a field minimisation procedure along X , followed by Z , after a correction current had already been applied along Y . Field measurements during the scan of each coil are purposefully taken displaced from the expected minimum in order to avoid

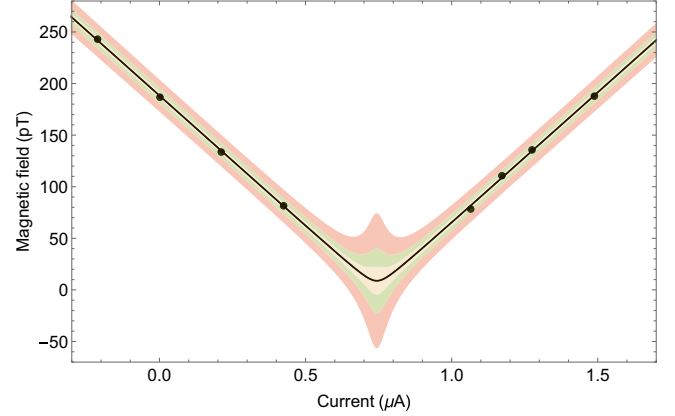


Fig. 6. A scan along Z fitted with equation (3) immediately after the minimum current from Figure 5 was applied.

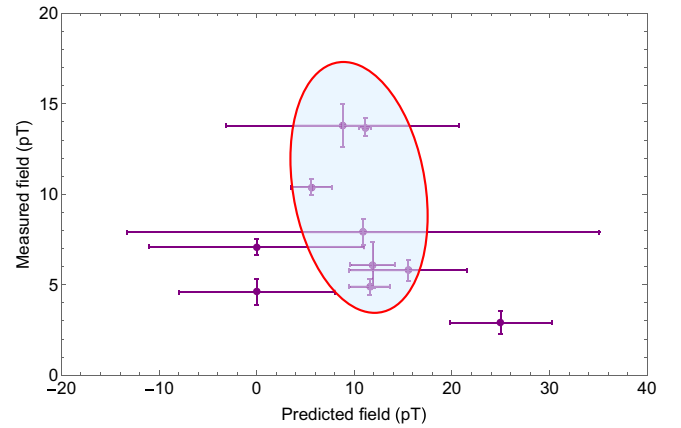


Fig. 7. Results of the minimisation measurements using the curve fit (Eq. (3)), along with the 1 σ region.

excessively long precession periods. Following the scan along X , the field minima was estimated to be 62(4) pT at a current of 1.57(2) μA . By scanning the current along Z immediately after applying the correction current in X , the magnitude of the 3-axis minimised field was predicted to be 9(12) pT with a final correction current of 0.74(2) μA . Both of the predicted field minima were calculated using equation (3) for the associated field component. Data taken after the current in the Z correction coil had been set to 0.74 μA determined the field to be 14(1) pT.

Ten such measurements were taken during the campaign and the results can be seen in Figure 7. From the plot, it can be seen that the absolute field attainable from carrying out the minimisation procedure using the hybrid magnetometer is of the order 10 pT, with a predicted field magnitude also around 10 pT.

After minimising the residual field through iteratively determining the three correction currents, it is possible to repeat the process to reach a lower residual field, reducing the coupling terms in equation (5). This process strongly depends on the magnetic field stability.

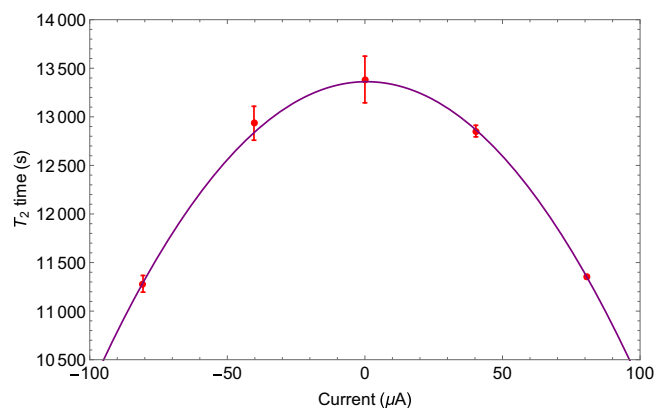


Fig. 8. Results of the T_2 measurement with peak T_2^* time of 13360(30) s.

5 Gradient minimisation

The effective T_2 decay time (T_2^*) of the ^3He precession signal is dependent on linear field gradients through the sample volume [22,23]. As the field gradient is changed through the measurement volume, by use of a set of gradient coils, the fit function

$$T_2^* = A(I - I_0)^2 + T_{2,\text{peak}}, \quad (6)$$

can be applied to determine the peak T_2^* time, $T_{2,\text{peak}}$, along with the correction current I_0 . To reduce the initial B_0 amplitude, the field was first minimised along two axes following the procedure set out in Section 4.1. The DAQ was instructed to apply five different currents to the X gradient coil, with two hours of data being collected at each value. The data from each current were analysed to extract the T_2^* time. The results of the measurement are shown in Figure 8. The results show that when the field gradient along X is minimised with an applied current of 0.6(4) μA , the peak T_2^* time has a value of 13360(30) s. If the optimum T_2 time is known, this value can be used to set limits on the remaining field gradients.

6 Conclusion

We have demonstrated the process of using ^{87}Rb zero-field magnetometers to probe the nearby precession of polarised ^3He , determining absolute low-field readings below 5 pT. It was shown that the minimisation procedure utilised in the 5-layer test magnetic shield at PSI results in a consistently achievable internal magnetic field at the 10 pT level. The results, a factor of 30 above the sensitivity limit from the Earth's rotation, were generally limited by the field stability inside the magnetic shielding. The results could further be improved by the use of an external field compensation system [24] around the shielding. Further measurements demonstrated the ability to also minimise magnetic field gradients through the volume of the ^3He cell. By providing absolute field readings in the pT range, the results presented here offer a new method for optimising the magnetic field environment in

a degaussed shield, by measuring residual magnetic fields and determining the relevant correction coil currents to cancel them. This then provides a near zero-field base onto which any desired magnetic field configuration can be applied using an appropriately designed field coil. We note that the ultimate stability and precision of the resulting field will also depend strongly on the current supplies used, and influence of the surrounding magnetic shielding. The $^3\text{He}/^{87}\text{Rb}$ hybrid magnetometer can then be used to also provide an absolute calibration for the applied field configuration if fields below 5 nT are used, or the system described in [15] could be used for larger fields. In the present work we are minimising the volume averaged residual field in a 70 mm diameter ^3He cell. Larger zero-field regions could in principle be accomplished if the position of the $^3\text{He}/^{87}\text{Rb}$ magnetometer was scannable within the shield, or by using an array of such sensors. This method can allow greatly improved understanding of the field configuration in shielded environments, providing a very useful tool for precision measurements.

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Availability of data and materials

The data used in this publication can be freely accessed at <https://doi.org/10.25377/sussex.8255852>

Author contribution statement

C. Abel and G. Bison conceived of the presented idea. C. Abel, G. Bison and W. C. Griffith. planned, collected and analysed the data. K. Kirch, B. Lauss and A. Mtchedlishvili conceived and constructed the magnetic shield. H. C. Koch and M. Pototschnig constructed the coil system. K. Kirch, B. Lauss and P. Schmidt-Wellenburg conceived and constructed the temperature stabilised environment. All authors discussed the results and contributed to the final manuscript with large inputs from C. Abel, G. Bison, W. C. Griffith, W. Heil, A. Schnabel, and J. Voigt.

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