

Investigation of the commensurate magnetic structure in the heavy-fermion compound CePt_2In_7 using magnetic resonant x-ray diffraction

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(Received 24 April 2017; revised manuscript received 19 July 2017; published 10 August 2017)

We investigated the magnetic structure of the heavy-fermion compound CePt_2In_7 below $T_N = 5.34(2)$ K using magnetic resonant x-ray diffraction at ambient pressure. The magnetic order is characterized by a commensurate propagation vector $\mathbf{k}_{1/2} = (\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$ with spins lying in the basal plane. Our measurements did not reveal the presence of an incommensurate order propagating along the high-symmetry directions in reciprocal space but cannot exclude other incommensurate modulations or weak scattering intensities. The observed commensurate order can be described equivalently by either a single- \mathbf{k} structure or by a multi- \mathbf{k} structure. Furthermore we explain how a commensurate-only ordering may explain the broad distribution of internal fields observed in nuclear quadrupolar resonance experiments [Sakai *et al.*, *Phys. Rev. B* **83**, 140408 (2011)] that was previously attributed to an incommensurate order. We also report powder x-ray diffraction showing that the crystallographic structure of CePt_2In_7 changes monotonically with pressure up to $P = 7.3$ GPa at room temperature. The determined bulk modulus $B_0 = 81.1(3)$ GPa is similar to those of the Ce-115 family. Broad diffraction peaks confirm the presence of pronounced strain in polycrystalline samples of CePt_2In_7 . We discuss how strain effects can lead to different electronic and magnetic properties between polycrystalline and single crystal samples.

DOI: [10.1103/PhysRevB.96.064414](https://doi.org/10.1103/PhysRevB.96.064414)

I. INTRODUCTION

Electrons can gain a large effective mass due to strong electronic correlations in crystals. Such materials are referred to as heavy-fermion compounds and often have complex phase diagrams due to the interplay of spin and electronic degrees of freedom. Of particular interest are the Ce-115 compounds CeMIn_5 ($M = \text{Co}, \text{Rh}, \text{Ir}$) that have been investigated for more than 15 years and yet their properties are still not completely understood [1]. These materials, which offer a unique playground to study quantum criticality [2], are part of the larger family $\text{Ce}_n\text{M}_m\text{In}_{3n+2m}$ ($M = \text{Co}, \text{Rh}, \text{Ir}, \text{Pd}, \text{Pt}$) derived from the simple cubic CeIn_3 : they are formed from CeIn_3 layers separated by $M\text{In}_2$ layers. This separation of the Ce planes makes them generally more two dimensional (2D) relative to the three-dimensional (3D) cubic CeIn_3 . Furthermore, the hybridization of the Ce $4f$ electrons with the conduction electron bands is controlled by the local environment of the In and M atoms [3,4]. It is therefore possible to investigate the effects of the dimensionality and the hybridization strength on the interplay between magnetism and superconductivity in these compounds.

CePt_2In_7 is a member of this family with $n = 1$ and $m = 2$. It is closely related to the Ce-115s and is obtained

by adding a second $M\text{In}_2$ plane in between the CeIn_3 planes. This larger separation of the planes containing Ce suggest that this system is more 2D than the Ce-115s. CePt_2In_7 crystallizes in a body-centered tetragonal structure with space group $I4/mmm$ and the magnetic Ce ion sits at the Wyckoff $2b$ positions [5,6]. It has an antiferromagnetic (AFM) order with $T_N \approx 5.5$ K at ambient pressure. This order is suppressed with pressure and a superconductivity dome emerges around the AFM quantum critical point (QCP), with a maximum $T_c = 2.1$ K near the critical pressure $P_c \approx 3.4$ GPa, and which is also where an effective mass enhancement is observed [7]. This phase diagram is very similar to the analogous compound CeRhIn_5 [8], which is often described as a two-dimensional analog of CeIn_3 . Quantum oscillations reveal that the microscopic electronic properties of CePt_2In_7 are more closely related to CeIn_3 than CeRhIn_5 , indicating that CePt_2In_7 is a better 2D analog of CeIn_3 [9]. The 2D nature of the electronic properties is also suggested by specific heat measurements [10]. Optical measurements indicate a hybridization strength in CePt_2In_7 similar to the one in CeIn_3 and CeRhIn_5 [11].

Nuclear quadrupolar resonance (NQR) measurements revealed the presence of two characteristic pressures in CePt_2In_7 [6]. The first one at $P^* = 2.4$ GPa corresponds to a transition from localized to itinerant Ce $4f$ electrons. The second one at $P_c \approx 3.4$ GPa corresponds to the AFM QCP. In CeRhIn_5 , these characteristic pressures are very close to each other and it was suggested that the superconductivity emerges from the Kondo breakdown QCP [12]. Indeed, recent theoretical work proposes an enhancement of singlet

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superconductivity near a Kondo breakdown QCP [13], which may explain the behavior of CePt₂In₇ and CeRhIn₅ [6,12]. The detailed understanding of CePt₂In₇ also requires an accurate description of its magnetic order at ambient pressure and its evolution (or stability) under pressure. However, up until now only limited details of the nature of the magnetic order have been reported. NQR measurements on polycrystalline samples indicate a commensurate order and suggest a propagation vector $(\frac{1}{2}, \frac{1}{2})$ in the basal plane [14]. On the other hand, the results obtained using the same technique applied to single crystals were interpreted in terms of a coexistence of commensurate and incommensurate orders [15]. From muon spin rotation measurements, a commensurate order was proposed for polycrystalline samples [16]. A possible reason for these discrepancies is that the inherently larger surface strain of grains in polycrystalline samples provides a means to enhance the stability of the commensurate order [15]. It was also observed that the superconducting dome is broader for powders than for single crystals, suggesting a commensurate order to be more favorable for superconductivity [17]. However, both direct measurements of the magnetic order and its propagation, and evidence for the proposed crystallographic strain in powder samples are yet to be reported.

Neutron scattering could clarify the bulk magnetic structure but it is challenging for CePt₂In₇ because of the generally small size of single crystals, the large neutron absorption cross section by In, and the small expected moment size. These limitations can be overcome by using magnetic resonant x-ray diffraction (MRXD) as an alternative scattering technique for determining the magnetic structure. We performed MRXD measurements on CePt₂In₇ and we report here a model for the magnetic order at $T = 1.8$ K and ambient pressure. We also report the pressure dependence of its crystallographic structure at room temperature up to $P = 7.3$ GPa, which changes monotonically in the range of applied pressure.

II. EXPERIMENTAL DETAILS

High-purity single crystals of CePt₂In₇ were synthesized as described previously [18]. The 0.38 mg sample used for the MRXD experiment was characterized by specific heat and magnetic susceptibility using a Quantum Design physical properties measurement system and magnetic property measurement system, respectively. The results are in good agreement with the previously reported measurements [18]. The long-range magnetic order is observed from a sharp peak in the specific heat at $T_N = 5.36(2)$ K [Fig. 2(c)] and the high purity of the sample is indicated by the absence of other peaks, compared to previous reports [18,19]. For the MRXD experiment, the platelike sample with the c axis perpendicular to the plate was fixed on a copper holder with silver ElectroDag 1415 and mounted in a Joule Thomson cryostat on the bending magnet XMaS beamline, at the European Synchrotron Radiation Facility. The measurements were carried out using a Vortex Si Drift Diode detector. The (220) reflection of a LiF analyzer crystal was used for the polarization analysis measurements. Except for photon energy dependent scans, all the measurements were carried out at $E = 6.166$ keV,

the Ce- L_{II} absorption edge. The azimuthal scans presented in Fig. 3 were corrected for x-ray absorption. The absorption correction was calculated by a finite element analysis assuming an absorption coefficient $\mu = 436.425$ mm⁻¹ for CePt₂In₇, a beam size of 0.7×0.3 mm², and a sample size of $0.79 \times 0.62 \times 0.02$ mm³. The accuracy of this correction for the magnetic peaks was verified by comparison with azimuthal scans measured on structural peaks.

Powder x-ray diffraction measurements under hydrostatic pressure were performed at the MS-X04SA beamline, Swiss Light Source at the Paul Scherrer Institut [20]. A 2D Pilatus 6M detector was used. LaB₆ was used as a standard for calibration of the detector position as well as the instrumental parameters. Single crystals of CePt₂In₇ were finely ground, mixed with quartz powder and loaded in a diamond-anvil pressure cell using methanol:ethanol 4:1 as a pressure medium. Quartz was used as an *in situ* pressure calibrant [21]. Measurements were performed with a photon wavelength $\lambda = 0.56491$ Å in the angular range $1^\circ < 2\theta < 35^\circ$ at room temperature ($T = 293$ K) up to a maximal pressure $P = 7.3$ GPa. The data reduction was performed with the DIOPTAS software [22] and FULLPROF was used for Rietveld refinement of the one-dimensional diffraction patterns [23].

III. EXPERIMENTAL RESULTS

A. Magnetic resonant x-ray diffraction

The magnetic order of CePt₂In₇ was successfully observed using MRXD, revealing unambiguously its commensurate propagation vector. Bragg peaks consistent with a propagation vector $\mathbf{k}_{1/2} = (\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$ were observed at $T = 1.8$ K. The magnetic origin of these Bragg peaks was verified by the resonance at the Ce- L_{II} absorption edge as well as polarization analysis. \mathbf{Q} scans around the magnetic Bragg peak $\mathbf{Q} = (0.5, -0.5, 6.5)$ are presented in Figs. 1(a)–1(c) at $T = 1.8$ K and can be compared with background scans done at 10 K. This magnetic Bragg peak has the same widths and shapes along H , K , and L as the structural Bragg peak $\mathbf{Q} = (1, -1, 6)$. This indicates that the magnetic peak widths are limited by the crystal mosaicity and that a 3D long-range magnetic order is achieved. Several other peaks consistent with $\mathbf{k}_{1/2}$ were measured. It was observed that all experimentally accessible magnetic Bragg peaks have nonzero intensity, indicating the absence of any selection rules of the magnetic structure.

The fluorescence intensity of the sample was measured as a function of the incident photon energy. It shows a maximum around $E_i = 6.167$ keV corresponding to the Ce- L_{II} absorption edge [Fig. 2(a)]. The intensity of the magnetic Bragg peak $\mathbf{Q} = (0.5, -0.5, 6.5)$ is strongly enhanced around this edge, indicating a resonant magnetic effect [24]. In contrast, the intensity of the structural Bragg peak $\mathbf{Q} = (1, -1, 6)$ shows a dip near this edge due to a larger absorption cross section. The magnetic nature of the Bragg peak $\mathbf{Q} = (0.5, -0.5, 6.5)$ is further confirmed by the polarization analysis. The polarization σ is defined to be perpendicular to the scattering plane and the polarization π is parallel to it [24]. In the electric dipole approximation of MRXD, charge scattering, related to the crystallographic structure, is allowed

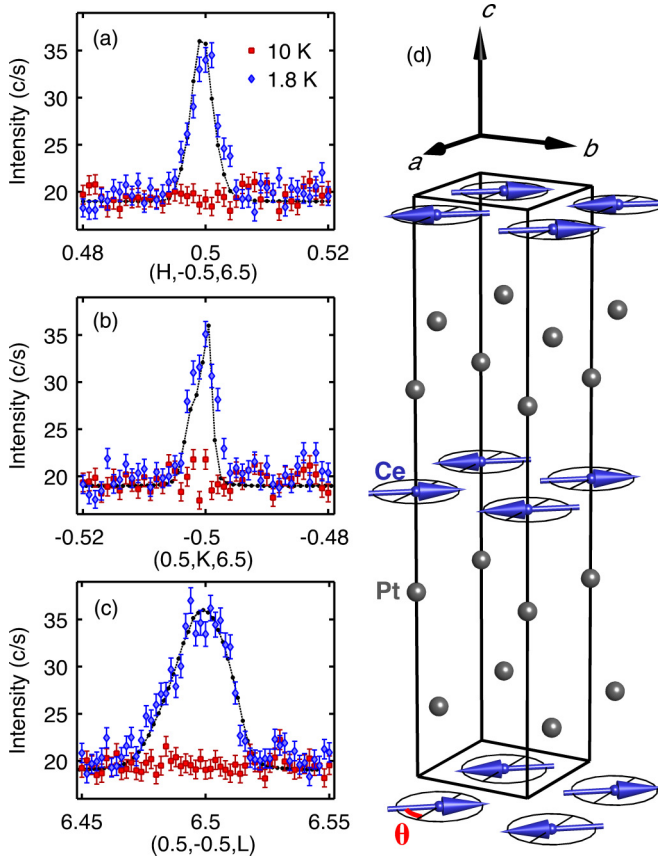


FIG. 1. (a) H scans, (b) K scans, and (c) L scans around the magnetic Bragg peak $\mathbf{Q} = (0.5, -0.5, 6.5)$ in the ordered state ($T = 1.8$ K) and in the paramagnetic state ($T = 10$ K). The connected black dots are corresponding scans around the structural Bragg peak $\mathbf{Q} = (1, -1, 6)$, and scaled to provide a comparison between the peak widths and shapes. (d) Single- \mathbf{k} magnetic structure of CePt_2In_7 with $\mathbf{k}_{1/2} = (\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$ and moments aligned in the basal plane at an angle θ from the a axis. Indium atoms have been omitted for clarity.

in the $\sigma - \sigma'$ channel and is forbidden in the $\sigma - \pi'$ channel. Magnetic scattering has the opposite behavior and appears in the $\sigma - \pi'$ channel and not in the $\sigma - \sigma'$ one [24]. The Bragg peak $\mathbf{Q} = (0.5, -0.5, 6.5)$ is present in the $\sigma - \pi'$ channel and absent in the $\sigma - \sigma'$ channel, clearly showing its magnetic nature [Fig. 2(b)]. This observation combined with the peak resonance at the Ce- L_{II} edge establish unambiguously the magnetic origin of the Bragg peaks with the propagation vector $\mathbf{k}_{1/2} = (\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$.

The temperature dependence of the magnetic Bragg peak $\mathbf{Q} = (0.5, -0.5, 6.5)$ has been measured from $T = 1.8$ K up to 7 K in the $\sigma - \pi'$ channel. The width and position of this peak are temperature independent from $T = 1.8$ K to T_N . The integrated intensity indicates a Néel temperature of $T_N = 5.34(2)$ K, as determined by a power-law fit above 4.4 K [Fig. 2(c)]. This transition temperature is in good agreement with the sharp peak observed in specific heat. The obtained critical exponent $\beta = 0.31(4)$ corresponds to a 3D Ising model with $\beta = 0.326$ or a 3D XY model with $\beta = 0.345$ [25]. Note that a beam injection occurred during the measurements at $T = 4.1$ K and that the intensity above and below this temperature cannot be compared accurately. However, a previous temperature de-

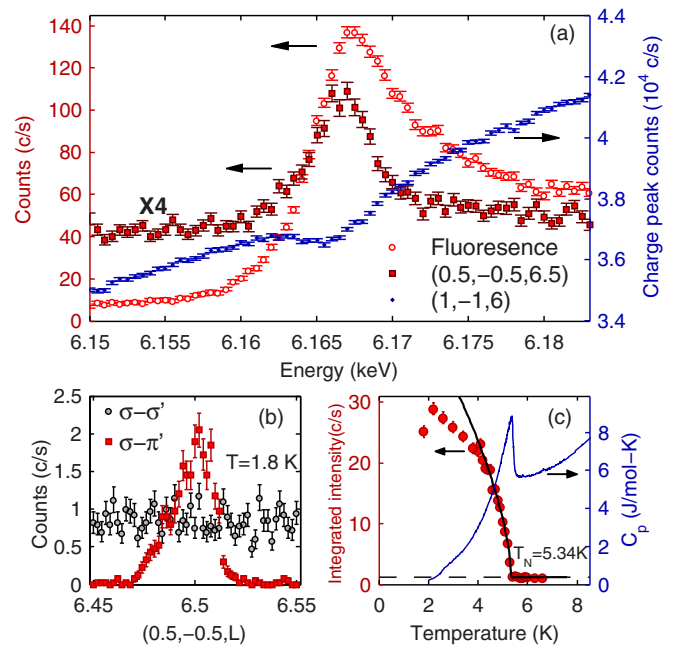


FIG. 2. (a) Photon energy scans near the L_{II} absorption edge of Ce: fluorescence of the sample, resonance of the magnetic Bragg peak $\mathbf{Q} = (0.5, -0.5, 6.5)$ at the absorption edge, and dip of the structural Bragg peak $\mathbf{Q} = (1, -1, 6)$. (b) L scan around the Bragg peak $\mathbf{Q} = (0.5, -0.5, 6.5)$ with polarization analysis, showing that all the signal is in the $\sigma - \pi'$ channel and therefore magnetic. The finite signal in the $\sigma - \sigma'$ channel is due to nonmagnetic background. (c) Temperature dependence of the integrated intensity of the magnetic Bragg peak $\mathbf{Q} = (0.5, -0.5, 6.5)$, showing the transition at $T_N = 5.34(2)$ K in agreement with the sharp peak in specific heat. The black line is a power-law fit to extract T_N . A beam injection occurred at $T = 4.1$ K and the intensity below and above this value cannot be compared directly.

pendence of the Bragg peak $\mathbf{Q} = (0.5, -0.5, 6.5)$ without the polarization analysis (not shown) does not have any feature at $T \approx 4$ K.

The magnetic structure of the propagation vector $\mathbf{k}_{1/2} = (\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$ was determined with the help of representation analysis performed with BASIREPS. [23]. Only two irreducible representations with nonzero basis functions are possible at the Ce position $(0, 0, 0.5)$ in the space group $I4/mmm$. There is Γ_1 , a two-dimensional irreducible representation with basis vector $(M_x, M_y, 0)$, and Γ_2 , a one-dimensional irreducible representation with basis vector $(0, 0, M_z)$. Both representations do not have selection rules, in agreement with our observations, and hence cannot be distinguished in this way.

We have determined that the structure must be described by Γ_1 with moments in the ab plane by performing azimuthal scans. These scans measure the intensity variation when the sample is rotated by azimuthal angle Ψ around the scattering vector \mathbf{Q} . In MRXD, the scattering intensity is proportional to $|\mathbf{F}(\mathbf{Q}) \cdot \mathbf{k}_f|^2$ where $\mathbf{F}(\mathbf{Q})$ is the magnetic structure factor and \mathbf{k}_f is the scattered photon wave vector [24]. Azimuthal rotations change the moment direction, modifying $\mathbf{F}(\mathbf{Q})$ relative to a fixed \mathbf{k}_f . The scattered intensity is therefore expected to change with Ψ and this can be compared with

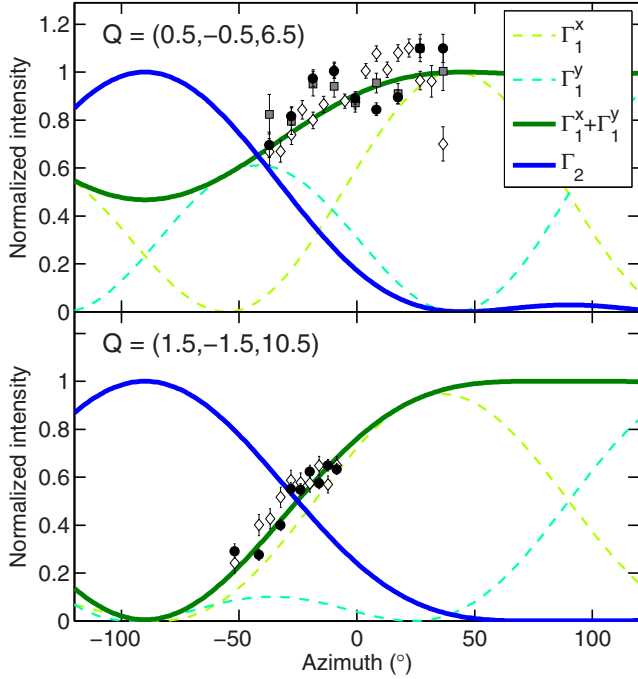


FIG. 3. Azimuthal scans on the magnetic Bragg peaks $\mathbf{Q} = (0.5, -0.5, 6.5)$ and $\mathbf{Q} = (1.5, -1.5, 10.5)$. The integrated intensities were corrected for absorption and scaled to compare to the models. Γ_1^x and Γ_1^y represent two domains with moments along a and b , respectively. Assuming equal population of both domains, the model agrees well with the results. Γ_2 represent the magnetic structure with moments along the c axis. Multiple datasets are represented by different symbols (see text).

that expected according to a magnetic structure model. The azimuth Ψ is defined relative to a reference Bragg peak, here chosen to be $\mathbf{Q} = (-1, -1, 0)$. The azimuthal angle is defined to be zero when the reference Bragg peak is in the scattering plane and forms the smallest angle with the incident photon wave vector \mathbf{k}_i .

Multiple datasets of azimuthal scans were collected and are represented by different symbols in Fig. 3. These datasets were collected in similar conditions (with and without optimizing the different rotation and translation motors) and all show the same general tendency. For both irreducible representations, the magnetic structure is collinear and the azimuthal scans correlate directly to the moment orientation. The theoretical azimuthal dependence curves for moments pointing along the a axis (Γ_1^x), the b axis (Γ_1^y), and the c axis (Γ_2) are shown in Fig. 3 for the magnetic Bragg peaks $\mathbf{Q} = (0.5, -0.5, 6.5)$ and $\mathbf{Q} = (1.5, -1.5, 10.5)$. Experimental results are overlaid and show that the system can be described by the coexistence of Γ_1^x and Γ_1^y domains with equal population. Since the axes a and b are equivalent, one would indeed expect that both domains are present. In general, if a domain exists with a moment pointing in a direction \mathbf{e} within the ab plane, a domain with a moment pointing in a direction \mathbf{e}' perpendicular to \mathbf{e} in the ab plane is expected with an equal population. It can be shown that the azimuthal dependence of $\Gamma_1^e + \Gamma_1^{e'}$ for any \mathbf{e} in the ab plane is exactly the same as that of $\Gamma_1^x + \Gamma_1^y$. Therefore, our results indicate the moments are in the ab plane but do not

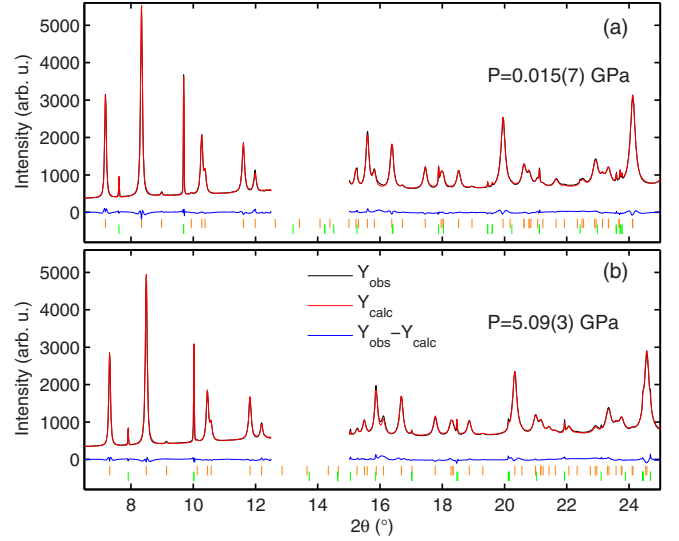


FIG. 4. X-ray diffraction pattern of CePt_2In_7 at $T = 293$ K for hydrostatic pressure of (a) $P = 0.015(7)$ GPa and (b) $P = 5.09(3)$ GPa. The region $12.5^\circ < 2\theta < 15^\circ$ has been excluded from the refinement. Orange and green tick marks indicate the Bragg peak positions for CePt_2In_7 and quartz, respectively.

allow us to determine their exact orientation. The magnetic structure for antiferromagnetically ordered moments pointing in the basal plane at an angle θ from the a axis is schematized in Fig. 1(d).

From previous NQR experiments, it was claimed that at $T = 1.6$ K an incommensurate magnetic order coexists with the commensurate order, and that the volume fraction of commensurate:incommensurate order was 0.25:0.75. In addition the maximal internal field due to the incommensurate order is determined to be slightly larger than the one from the commensurate order, suggesting a similar moment size for both orders. For these two reasons, the magnetic peak intensities originating from the incommensurate order can be expected to be similar to those of the commensurate order. However, no evidence for incommensurate magnetic peaks was found in our MRXD experiment from scans along the high-symmetry directions in reciprocal space. Measurements were carried out at $T = 1.8$ K for $\mathbf{Q} = (0.5, -0.5, L)$ from $L = 6$ to 8, $\mathbf{Q} = (H, \bar{H}, 7)$ from $H = 0$ to 1.2, $\mathbf{Q} = (H, \bar{H}, 6.5)$ from $H = 0$ to 1.2, and $\mathbf{Q} = (H, 0, 6.5)$ from $H = 0$ to 1.5. This rules out likely incommensurate propagation vectors similar to those of other incommensurate magnetic phases in Ce-based heavy-fermion compounds [26–31], but we cannot exclude the presence of incommensurate modulations propagating elsewhere in reciprocal space.

B. Powder x-ray diffraction under pressure

Powder x-ray diffraction patterns of CePt_2In_7 are shown in Fig. 4 at hydrostatic pressures $P = 0.015(7)$ GPa and $P = 5.09(3)$ GPa for a representative 2θ angular range. The general crystallographic structure, previously reported by Klimczuk *et al.* [5], was confirmed by Rietveld refinement using FULLPROF [23]. Two strong diffraction peaks from CePt_2In_7

appear in the angular range $12.5^\circ < 2\theta < 15^\circ$ and this region has been excluded from the refinement to improve the sensitivity of the fit to weak features over the full angular range. Importantly, the diffraction peak profiles due to the CePt_2In_7 sample are significantly broader than the instrumental resolution and this can be attributed to strain. The presence of strain in polycrystalline samples of CePt_2In_7 was inferred previously in NQR measurements [15]. Our measurements remained in a hydrostatic regime up to the maximal applied pressure, as confirmed by the pressure independent widths of peaks due to scattering from quartz. However, the peak widths of CePt_2In_7 gradually broadened above $P \approx 5$ GPa, which show a loss of the structural integrity in terms of either a larger strain or breaking of crystallites into smaller particles.

The refinement of the diffraction patterns was performed sequentially for increasing pressure and the results are presented in Fig. 5. We observe no obvious changes of the crystallographic structure related to the characteristic pressures $P^* = 2.4$ GPa and $P_c \approx 3.4$ GPa. The lattice constants a and c change monotonically up to the maximal applied pressure $P = 7.3$ GPa [Figs. 5(a)–5(c)]. The Birch-Murnaghan equation of state was used to relate the crystal volume V to the applied pressure P :

$$P(V) = \frac{3}{2}B_0(v^7 - v^5)\left[1 - \frac{3}{4}(4 - B'_0)(v^3 - 1)\right], \quad (1)$$

where B_0 , and B'_0 are, respectively, the initial bulk modulus and its derivative, and $v = (V_0/V)^{1/3}$ [32]. By fitting this equation to the data shown in Fig. 5(d), we obtain $B_0 = 81.1 \pm 0.3$ GPa and $B'_0 = 5.8 \pm 0.1$. Using the simple Murnaghan equation [33] results in the same fitted values for B_0 and B'_0 within errors. These values are similar to those reported for other members of the $\text{Ce}_n\text{M}_m\text{In}_{3n+2m}$ family [34]. In these compounds, it was observed that adding $M\text{In}_2$ layers stiffens the structure and increases the bulk modulus: $B_0 = 67$ GPa for CeIn_3 , average $B_0 = 70.4$ GPa for Ce_2MIn_8 (two layers CeIn_3 + one layer $M\text{In}_2$), and average $B_0 = 81.4$ GPa for CeMIn_5 (one layer CeIn_3 + one layer $M\text{In}_2$). The addition of a second $M\text{In}_2$ layer in CePt_2In_7 relative to CeMIn_5 could then be expected to stiffen the lattice further. However, the bulk moduli appear very similar for CePt_2In_7 and the Ce-115s .

In CePt_2In_7 , the Ce and Pt atoms sit at Wyckoff positions $2b$ and $4e$, respectively, and the In atoms are distributed on three different positions [In(1) at $2a$, In(2) at $4d$ and In(3) at $8g$]. The only adjustable fractional coordinates in the structure of CePt_2In_7 are the Z positions of the Pt and In(3) atoms. The fractional coordinate Z of In(3) changes monotonically with pressure [Fig. 5(e)] and the one of Pt is pressure independent [Fig. 5(f)]. This indicates a nonuniform compression along the c axis, with the strongest contraction occurring between the In(3) planes and the Ce-In(1) planes [see Fig. 5(g)]. The pressure dependence of various bond lengths is presented in Figs. 5(h)–5(j) and they all decrease monotonically with increasing pressure. Interestingly, the Ce-In(3) bond is more significantly affected by pressure than the Ce-In(1) bond [Fig. 5(h)]. Since the Ce-In coupling is expected to be the strongest with the out-of-plane In(3) atoms [4,35], this change in distortion around the Ce atoms could modify significantly the ground-state Ce wave function [3].

IV. DISCUSSION

As mentioned previously, the pressure-temperature phase diagram of CeRhIn_5 is very similar to that of CePt_2In_7 . Their magnetic structures at ambient pressure also share similarities: both have an antiferromagnetic order in the basal plane with moments lying in that plane [26]. However, the ordering in CeRhIn_5 is incommensurate along the c axis in contrast with the commensurate ordering in CePt_2In_7 . While CeCoIn_5 and CeIrIn_5 do not order magnetically at ambient pressure and zero magnetic field, it is possible to induce magnetic order with doping. In particular, substituting the Co or Ir sites with Rh leads to the coexistence of an incommensurate order with $\mathbf{k} = (\frac{1}{2}, \frac{1}{2}, \delta)$ and a commensurate order with $\mathbf{k} = (\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$ for a range of doping values [27–30]. It was shown for $\text{CeRh}_{0.7}\text{Ir}_{0.3}\text{In}_5$ specifically that the moments lie in the basal plane for both the commensurate and incommensurate orders. Doping the In site with Cd in CeCoIn_5 also stabilizes a commensurate order with $\mathbf{k} = (\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$ [36]. On the other hand, substituting Ce by Nd in CeCoIn_5 leads to a propagation vector $\mathbf{k} = (\frac{1}{2} - \delta, \frac{1}{2} - \delta, \frac{1}{2})$ with $\delta = 0.05$ [31], suggesting a spin-density wave in the basal plane with fundamentally different properties from the localized moment magnetism in CeRhIn_5 and CePt_2In_7 .

In these systems, superconductivity emerges in the vicinity of an AFM QCP, suggesting a magnetically driven pairing mechanism of superconductivity. The knowledge of the magnetic structure is therefore a crucial element for identifying the magnetic fluctuations responsible for this electron-electron coupling. The AFM order $(\frac{1}{2}, \frac{1}{2})$ in the basal plane prevails in these systems and CePt_2In_7 appears as a new example where magnetic fluctuations associated with this AFM order are the pairing glue of the pressure-induced superconductivity. It is important to note that the magnetic structure of CePt_2In_7 might change under pressure but it is unlikely to change the order in the basal plane. For example, the propagation vector in CeRhIn_5 changes under pressure but the order in the basal plane is conserved [37–40].

Based on NQR experiments, it was suggested that in single crystals of CePt_2In_7 there is a coexistence of commensurate and incommensurate orders at ambient pressure [6,15]. Specifically, sharp peaks in the spectrum can be attributed to a basal plane AFM order with moments pointing along the a axis or the b axis. This was interpreted as a commensurate order. On the other hand, broad features are also observed in the spectrum and were attributed to a distribution of internal fields at the In(2) and In(3) sites. This was interpreted as an incommensurate order similar to that of CeRhIn_5 [41].

Our results presented in Sec. III A confirm the presence of a commensurate order but do not reveal the presence of an incommensurate order along the high-symmetry directions in reciprocal space indicated in Sec. III A. The scenario involving the coexistence of both commensurate and incommensurate orders remains a possibility: we cannot rule out incommensurate modulations propagating elsewhere in reciprocal space, and the volume fraction and/or moment size could be too small to be detected under our current experimental conditions.

On the other hand, we propose an alternative interpretation of the broad features observed in the NQR experiments that

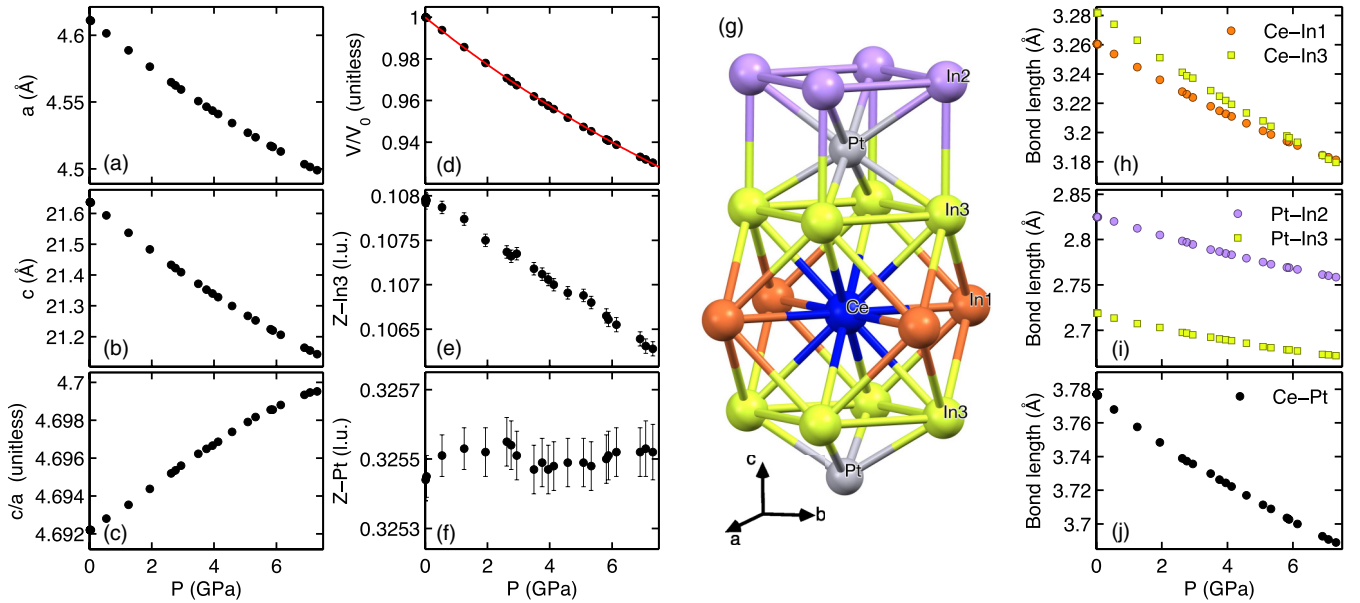


FIG. 5. Pressure dependence of the crystallographic structure at $T = 293$ K. (a)–(c) Lattice constants a and c and ratio c/a as a function of pressure. (d) Unit cell volume as a function of pressure, with the fit to the Birch-Murnaghan equation of state represented by the red solid line. (e),(f) Pressure dependence of the fractional coordinate Z of Pt and In(3) atoms. (g) Crystallographic structure of CePt_2In_7 around the Ce and Pt atoms, where In atoms are colored by their site symmetry. (h)–(j) Pressure dependence of different bond lengths.

do not require the presence of an additional incommensurate order. With no restriction on the precise moment direction in the basal plane provided by our MXRD experiments, the distribution of internal fields observed by NQR could be generated if either the moment directions in the ab plane fluctuate, or there exist multiple domains with different moment orientations (different values of θ in Fig. 1). This commensurate-only scenario for the magnetic order in CePt_2In_7 requires a coexistence of domain types; those with arbitrary moment orientations in the ab plane as outlined above, and those where the moments are rigidly aligned with the a and b axes. Here crystal strain could play an important role in stabilizing one type of domain over the other.

In NQR experiments, different results for the reported spectra are obtained from polycrystalline and single crystal samples of CePt_2In_7 [14,15]. These discrepancies are readily attributable to crystal/surface strain effects that vary in propensity with the sample crystallite size. Indeed, this is supported by the broad structural peaks in our high-resolution powder x-ray diffraction experiment on CePt_2In_7 . In the NQR studies only sharp features are observed for powder samples, in contrast with the presence of broad features for single crystals. Furthermore, applied pressure on single crystals suppresses the contribution of the broad features [15]. Taken together, these two effects indicate that strain, either from surface strain from the grains in polycrystalline samples or stimulated by pressure, promotes the ordering with moments aligned along the a axis or the b axis. At the same time, in the absence of strain, the moments may align along an arbitrary direction in the ab plane. In this scenario, enhanced strain thus leads to an effective in-plane anisotropy that favors the alignment of the moments along the a axis or b axis.

It is interesting then to note that the superconductivity is stabilized in a wider pressure range in powder samples and that it only appears in single crystals when the NQR signature interpreted in terms of incommensurate order is completely suppressed [15,17]. This suggests that domains with moments not aligned along the a axis or the b axis are detrimental to the formation of superconductivity in CePt_2In_7 .

Finally, we note that even if the magnetic structure presented in Sec. III A is the simplest solution to explain the results, it is not the only possible one. Since the lattice of CePt_2In_7 is body centered, the propagation vector $\mathbf{k}_{1/2}$ is not equivalent to $-\mathbf{k}_{1/2}$. This can lead either to two different \mathbf{k} domains, which was assumed in Sec. III A, or a multi- \mathbf{k} structure, as observed, for example, in the heavy fermion CeRh_2Si_2 , which also has a body-centered tetragonal lattice [42]. A complete description of the multi- \mathbf{k} structure in CePt_2In_7 is given in the Appendix. In such a multi- \mathbf{k} structure, the moments between the nearest-neighboring Ce layers can be noncollinear, while all the moments are collinear in a single- \mathbf{k} structure. This noncollinearity suggests an effective decoupling of the nearest-neighbor layers while keeping a coupling to the next-nearest-neighbor planes, consequently forming two decoupled yet interpenetrating sublattices. This scenario is plausible for the body-centered tetragonal lattice because of the presence of competing interactions. It was even suggested theoretically that the frustration in body-centered tetragonal lattices can destabilize long-range magnetic order and lead to spin liquid states in heavy-fermion compounds [43]. The aforementioned discussion about the moment directions in the single- \mathbf{k} model, and its application for consistently explaining previously reported NQR spectra, can also be done using the multi- \mathbf{k} structure. Our results do not allow us to establish unambiguously if the single- \mathbf{k} structure or the multi- \mathbf{k} structure is the correct one. In fact, these two scenarios cannot be

distinguished in a simple scattering experiment; doing so would require the application of either uniaxial strain or magnetic fields to control the magnetic domain formation in a single crystal sample.

V. SUMMARY

We have shown that the crystallographic structure of CePt_2In_7 changes monotonically with pressure up to $P = 7.3$ GPa at room temperature. We also investigated the magnetic order of CePt_2In_7 at ambient pressure below $T_N = 5.34(2)$ K by magnetic resonant x-ray diffraction. This order is characterized by a commensurate propagation vector $\mathbf{k}_{1/2} = (\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$. The magnetic origin of these diffraction peaks was confirmed by their resonance at the Ce- L_{II} absorption edge and by polarization analysis. Azimuthal scans confirm that the moments lie in the basal plane. The magnetic structure can be described by a single- \mathbf{k} structure or by a multi- \mathbf{k} structure. Both structures cannot be distinguished in a simple scattering experiment as reported here and the single- \mathbf{k} structure is discussed for simplicity. The presence of incommensurate order in CePt_2In_7 was previously reported based on NQR experiments. Our measurements could not reveal the presence of such an order but are insufficient to exclude it completely. Using our results we propose a new scenario for the ambient pressure ground state of CePt_2In_7 that is described only by commensurate magnetic order; namely, a coexistence of domains wherein the moments are either rigidly aligned along the a and b axes, or arbitrarily aligned within the ab plane. Crystal strain is argued to be an effective tuning parameter for controlling the relative volume fractions of the two types of domain, thus providing a means for a consistent description of both the scattering data reported here, and previously reported NQR spectra obtained on both polycrystalline and single crystal samples.

Note added. Recently, we became aware of another report where the magnetic structure of CePt_2In_7 was investigated using neutron diffraction [44]. In agreement with our results, they report a commensurate propagation vector $\mathbf{k}_{1/2}$ and moments lying in the basal plane. The reported structure corresponds to a multi- \mathbf{k} structure with noncollinear moments.

ACKNOWLEDGMENTS

The authors are thankful to M. Kenzelmann, D. G. Mazzone, and F. Ronning for fruitful discussions. This research received support from the Natural Sciences and Engineering Research Council of Canada (Canada). Work at Los Alamos National Laboratory was performed under the auspices of the U.S. Department of Energy, Office of Basic Energy Sciences, Division of Materials Sciences and Engineering. XMaS is a UK midrange facility supported by EPSRC.

APPENDIX: SINGLE- \mathbf{k} AND MULTI- \mathbf{k} STRUCTURES

1. Single- \mathbf{k} structure

The simplest magnetic structure model for CePt_2In_7 is described by the single propagation vector $\mathbf{k}_{1/2}$. This is represented in Fig. 6(a). In the unit cell, there are two Ce ions which are related by the body-centering symmetry. They

are distinguished by the blue and green colors in Fig. 6. For a general single- \mathbf{k} structure, the moments \mathbf{m}_b and \mathbf{m}_g at the blue and green sites, respectively, are expressed as

$$\mathbf{m}_b = \begin{pmatrix} M \cos \theta \\ M \sin \theta \\ 0 \end{pmatrix} \cos \Phi, \quad (\text{A1})$$

$$\mathbf{m}_g = \begin{pmatrix} M \cos \theta \\ M \sin \theta \\ 0 \end{pmatrix} \cos \left(\Phi \pm \frac{\pi}{2} \right). \quad (\text{A2})$$

Here the parameter θ is the angle of the moment in the ab plane, which can take any value. To reproduce the data, the presence of two equally populated domains with θ and $\theta + 90^\circ$ is assumed. The parameter Φ is a global phase that cannot be measured with scattering techniques. For physical reasons, we chose $\Phi = \frac{\pi}{4}$ to generate equal moments for \mathbf{m}_b and \mathbf{m}_g . The single- \mathbf{k} structure is therefore defined by

$$\mathbf{m}_b = \frac{1}{\sqrt{2}} \begin{pmatrix} M \cos \theta \\ M \sin \theta \\ 0 \end{pmatrix}, \quad (\text{A3})$$

$$\mathbf{m}_g = -\frac{1}{\sqrt{2}} \begin{pmatrix} M \cos \theta \\ M \sin \theta \\ 0 \end{pmatrix}. \quad (\text{A4})$$

2. Multi- \mathbf{k} structure

Due to the body-centering symmetry, $+\mathbf{k}_{1/2}$ and $-\mathbf{k}_{1/2}$ are not equivalent and therefore, a magnetic structure can form that is composed of two propagation vectors. In a general way, the moments are defined at the blue and green sites, respectively, by

$$\mathbf{m}_b^{+\mathbf{k}} + \mathbf{m}_b^{-\mathbf{k}} = \begin{pmatrix} M \cos \theta_1 \\ M \sin \theta_1 \\ 0 \end{pmatrix} \cos \Phi_1 + \begin{pmatrix} M \cos \theta_2 \\ M \sin \theta_2 \\ 0 \end{pmatrix} \cos \Phi_2, \quad (\text{A5})$$

$$\mathbf{m}_g^{+\mathbf{k}} + \mathbf{m}_g^{-\mathbf{k}} = \begin{pmatrix} M \cos \theta_1 \\ M \sin \theta_1 \\ 0 \end{pmatrix} \cos \left(\Phi_1 + \frac{\pi}{2} \right) + \begin{pmatrix} M \cos \theta_2 \\ M \sin \theta_2 \\ 0 \end{pmatrix} \cos \left(\Phi_2 - \frac{\pi}{2} \right), \quad (\text{A6})$$

where θ_1 and Φ_1 are related to the propagation $+\mathbf{k}_{1/2}$, and θ_2 and Φ_2 are related to $-\mathbf{k}_{1/2}$. It is again assumed that there are two equally populated domains with $\{\theta_1, \theta_2\}$ and $\{\theta_1 + 90^\circ, \theta_2 + 90^\circ\}$. Experimentally, this gives exactly the same scattering as the single- \mathbf{k} structure. We must choose Φ_1 and Φ_2 to have equal moments on the blue and green sites for any θ_1 and θ_2 . An elegant choice is $\Phi_1 = \frac{n\pi}{2}$ and $\Phi_2 = \frac{(n+1)\pi}{2}$ where n is an integer. It evidences the decoupling of the nearest-neighbor layers. For $n = 0$, we obtain

$$\mathbf{m}_b^{+\mathbf{k}} + \mathbf{m}_b^{-\mathbf{k}} = \begin{pmatrix} M \cos \theta_1 \\ M \sin \theta_1 \\ 0 \end{pmatrix} \quad (\text{A7})$$

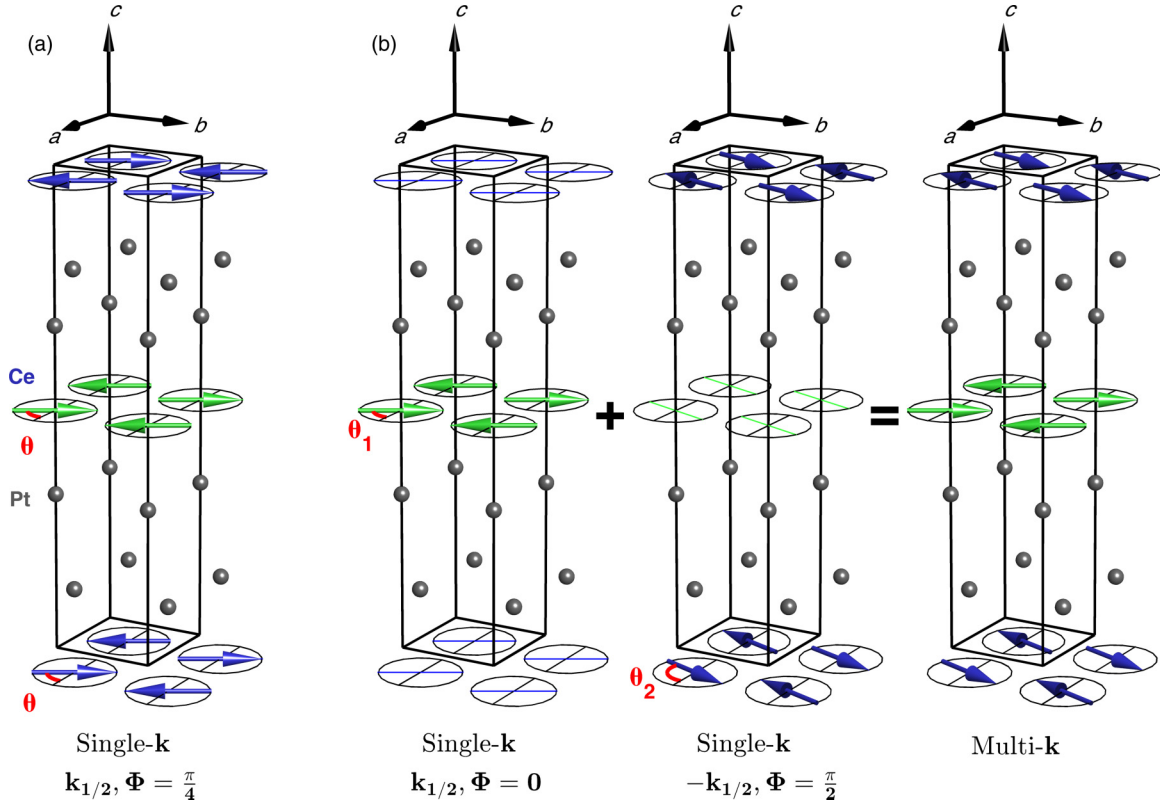


FIG. 6. (a) Single- \mathbf{k} structure with the propagation vector $\mathbf{k}_{1/2}$ and the phase factor $\Phi = \frac{\pi}{4}$. (b) Single- \mathbf{k} structures with $\mathbf{k} = \mathbf{k}_{1/2}$, $\Phi = 0$ and $\mathbf{k} = -\mathbf{k}_{1/2}$, $\Phi = \frac{\pi}{2}$, combined to form a multi- \mathbf{k} structure. The angles θ_1 and θ_2 indicate the moment direction in the basal plane relative to the a axis for each single- \mathbf{k} structure.

and

$$\mathbf{m}_g^{+\mathbf{k}} + \mathbf{m}_g^{-\mathbf{k}} = \begin{pmatrix} M \cos \theta_2 \\ M \sin \theta_2 \\ 0 \end{pmatrix}. \quad (\text{A8})$$

The structure is therefore defined by three parameters: the moment size M , the angle θ_1 of the first propagation vector, and the angle θ_2 of the second propagation vector. While M is expected to be constant, θ_1 and θ_2 can take any value. Note that the single- \mathbf{k} structure is obtained if $\theta = \theta_1 = \theta_2$.

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